

STATIC &

STATIONARY:

CAUSAL

CURIOSITIES &

FUNDAMENTAL

CYCLE

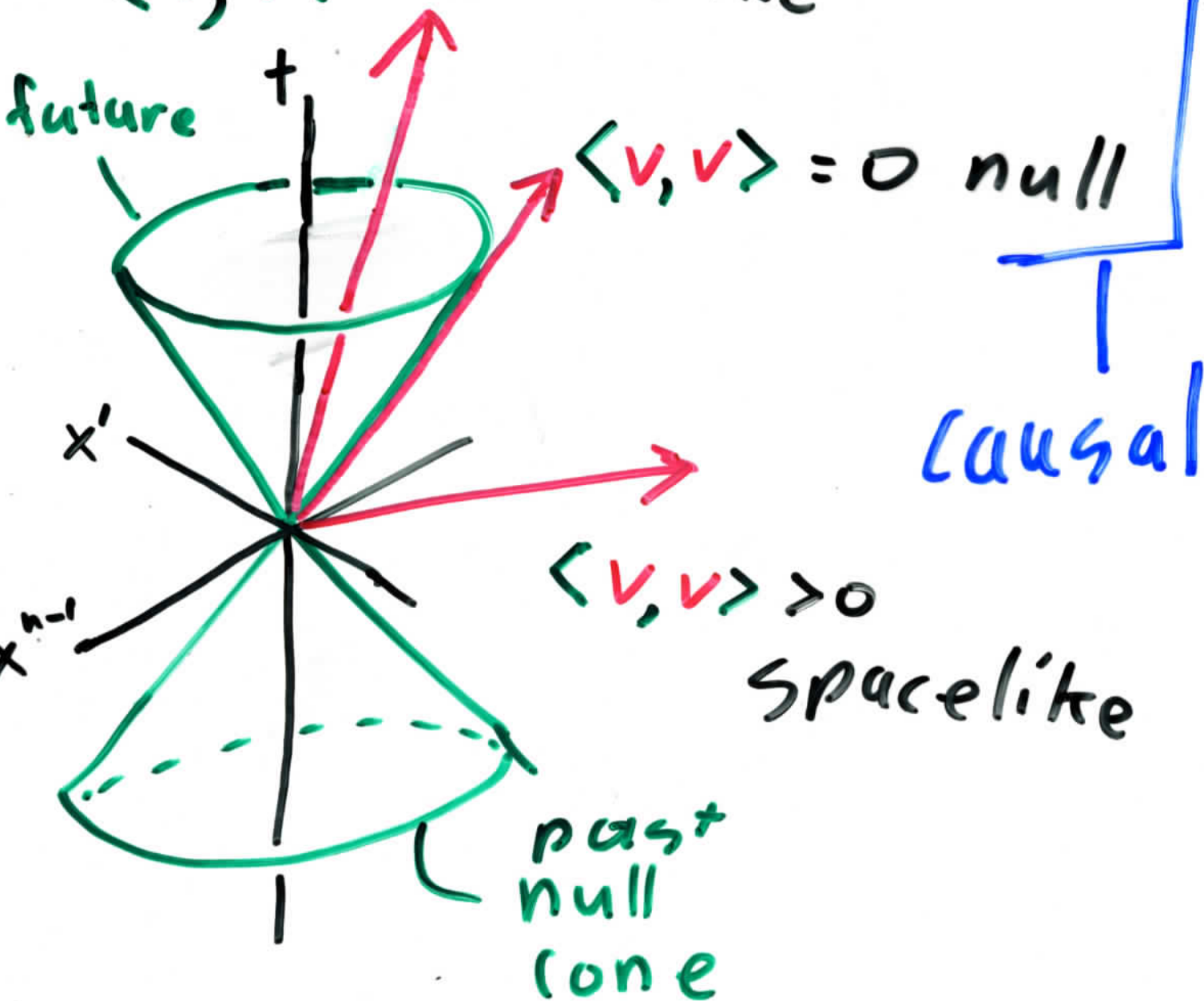
Stacey
Harris
Saint Louis
Univ.

Lorentz-signature metric on \mathbb{R}^n

$$-dt^2 + (dx^1)^2 + \dots + (dx^{n-1})^2$$

$$\langle v, w \rangle = -v^0 w^0 + v^1 w^1 + \dots + v^{n-1} w^{n-1}$$

$\langle v, v \rangle < 0$ timelike



$\langle v, v \rangle = 0$ null

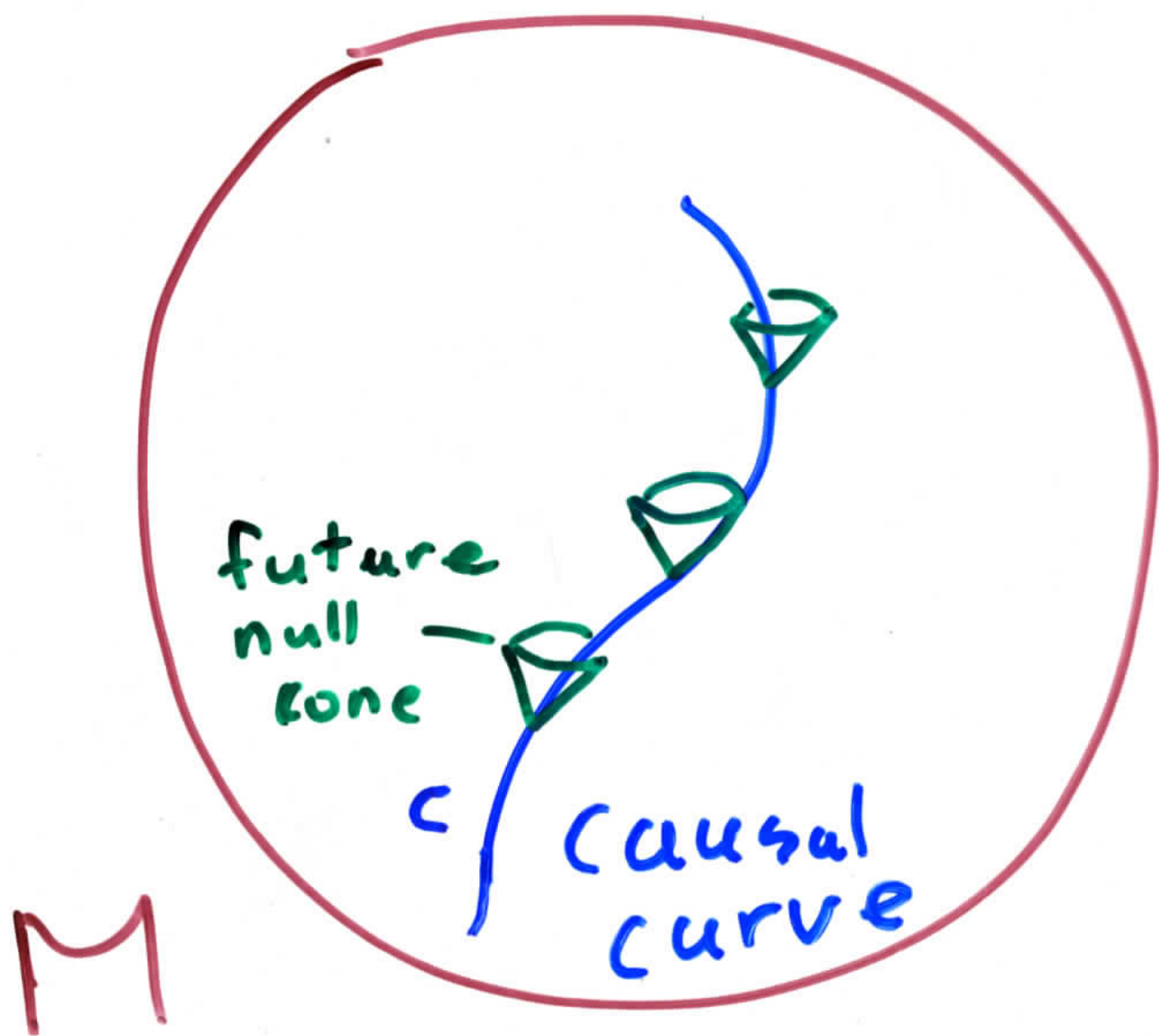
causal

$\langle v, v \rangle > 0$

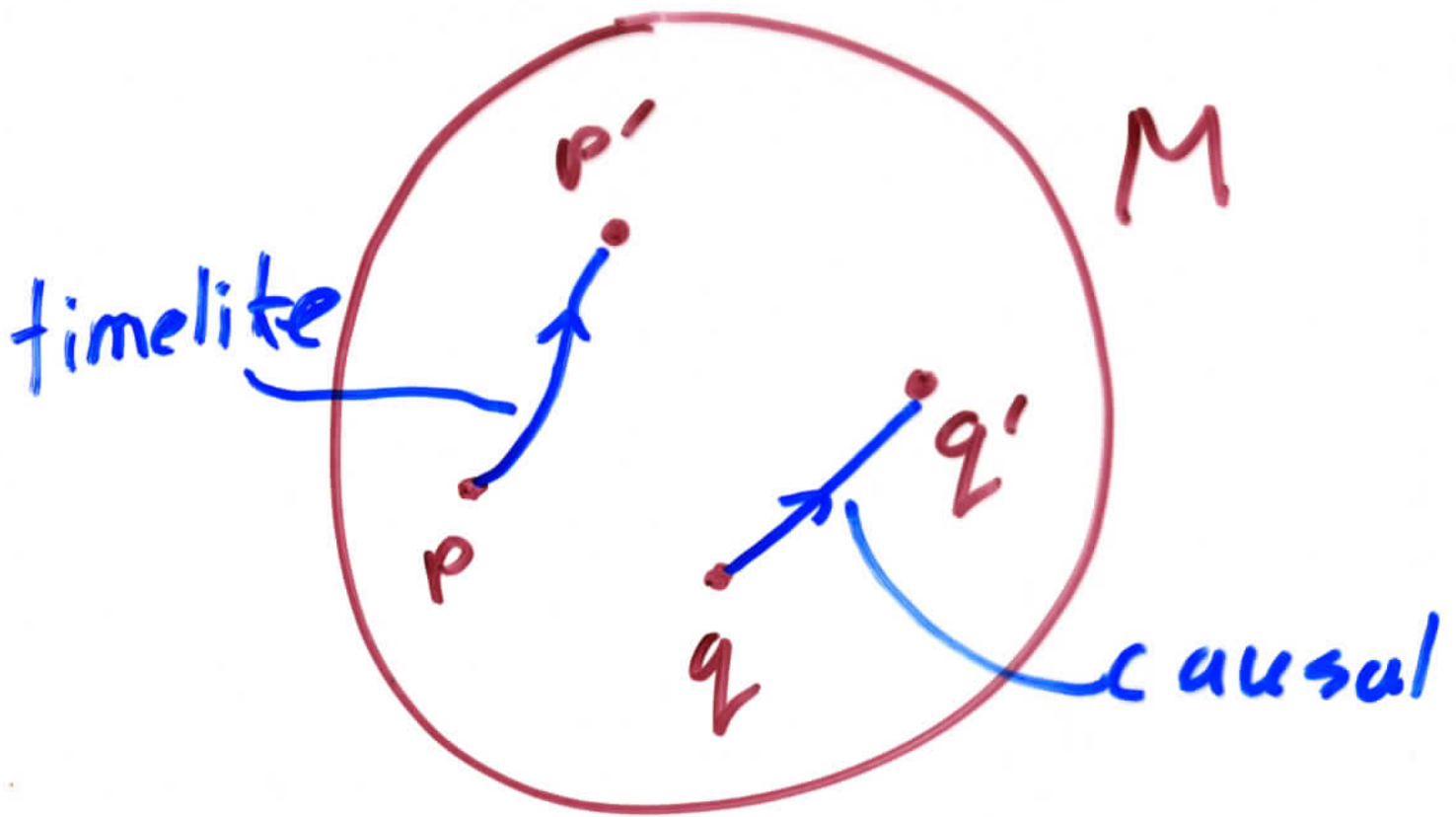
spacelike

past
null
cone

spacetime M^n
manifold w/ Lorentz
metric in each $T_p M$



$$\langle \dot{c}, \dot{c} \rangle \leq 0$$



$p \ll p'$ \exists future-timelike curve $p \rightarrow p'$
 |
 chronologically precedes
 (massive particles)

$q \leq q'$ \exists future-causal curve $q \rightarrow q'$
 |
 causally precedes
 (photons)

M is

chronological:

no closed timelike
curves



causal:

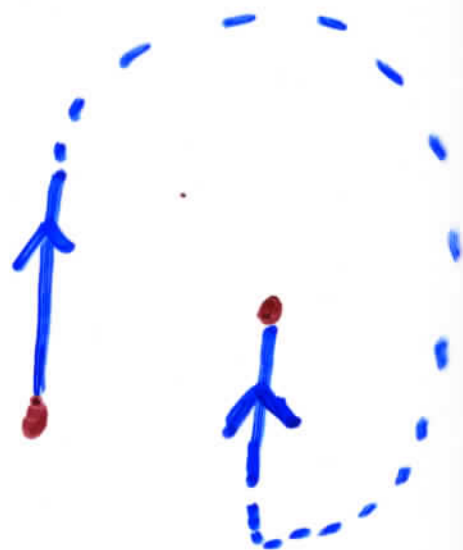
no closed causal
curves



strongly causal:

local causality
= global causality

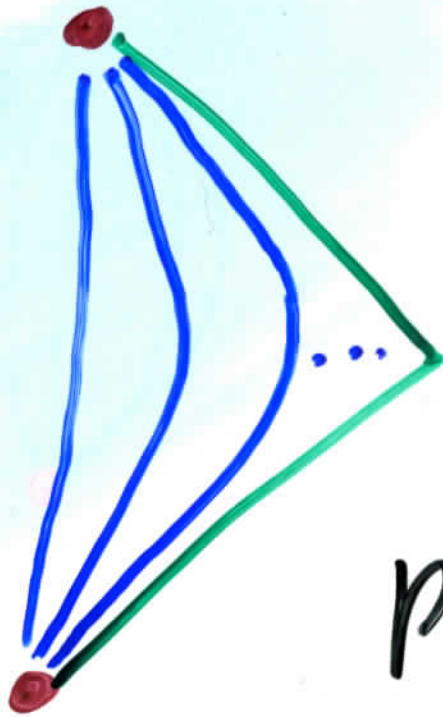
no



M is

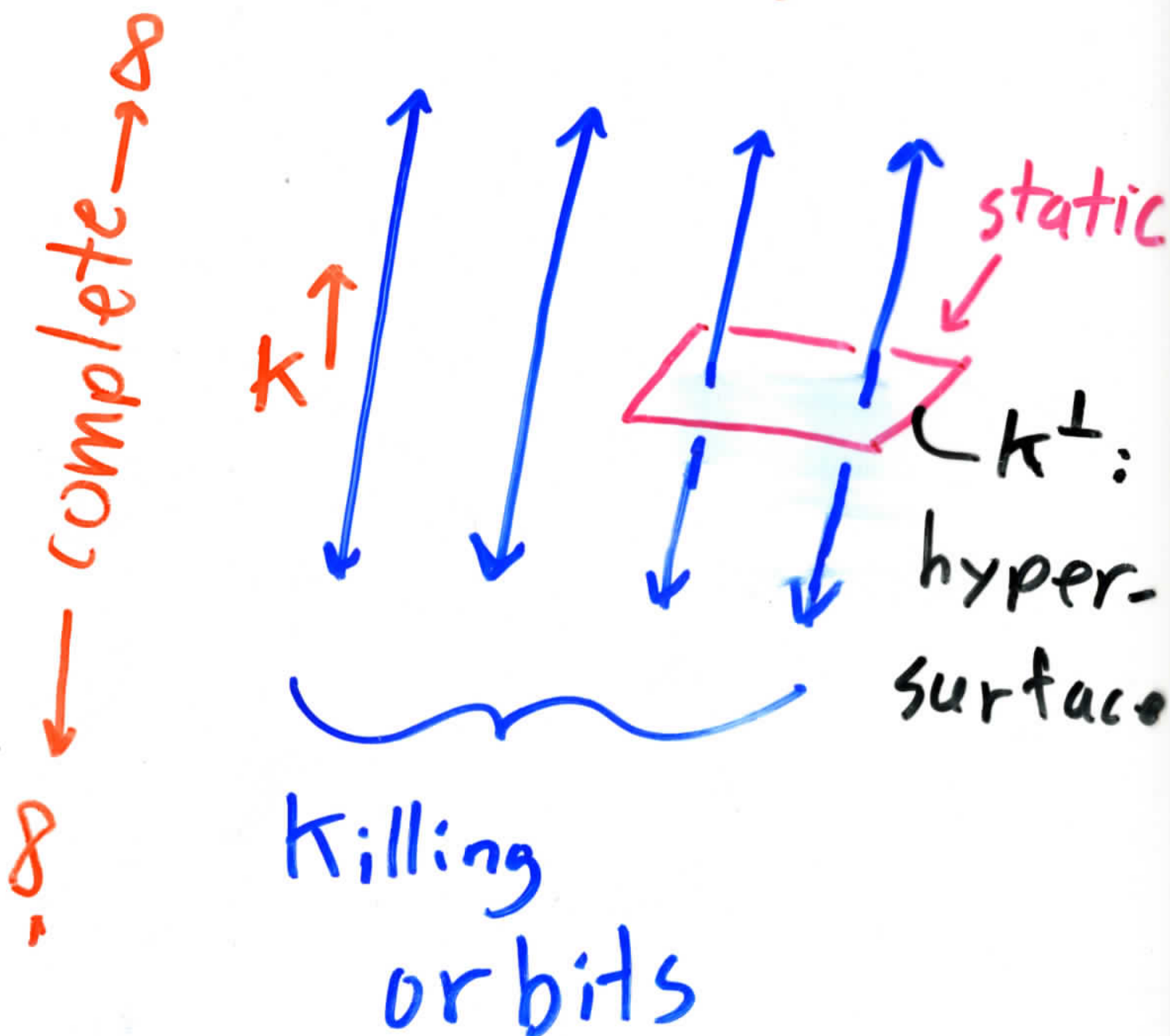
globally hyperbolic:

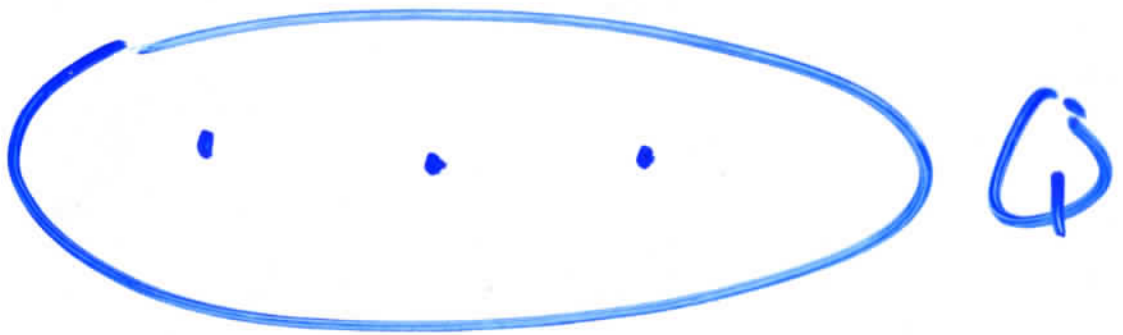
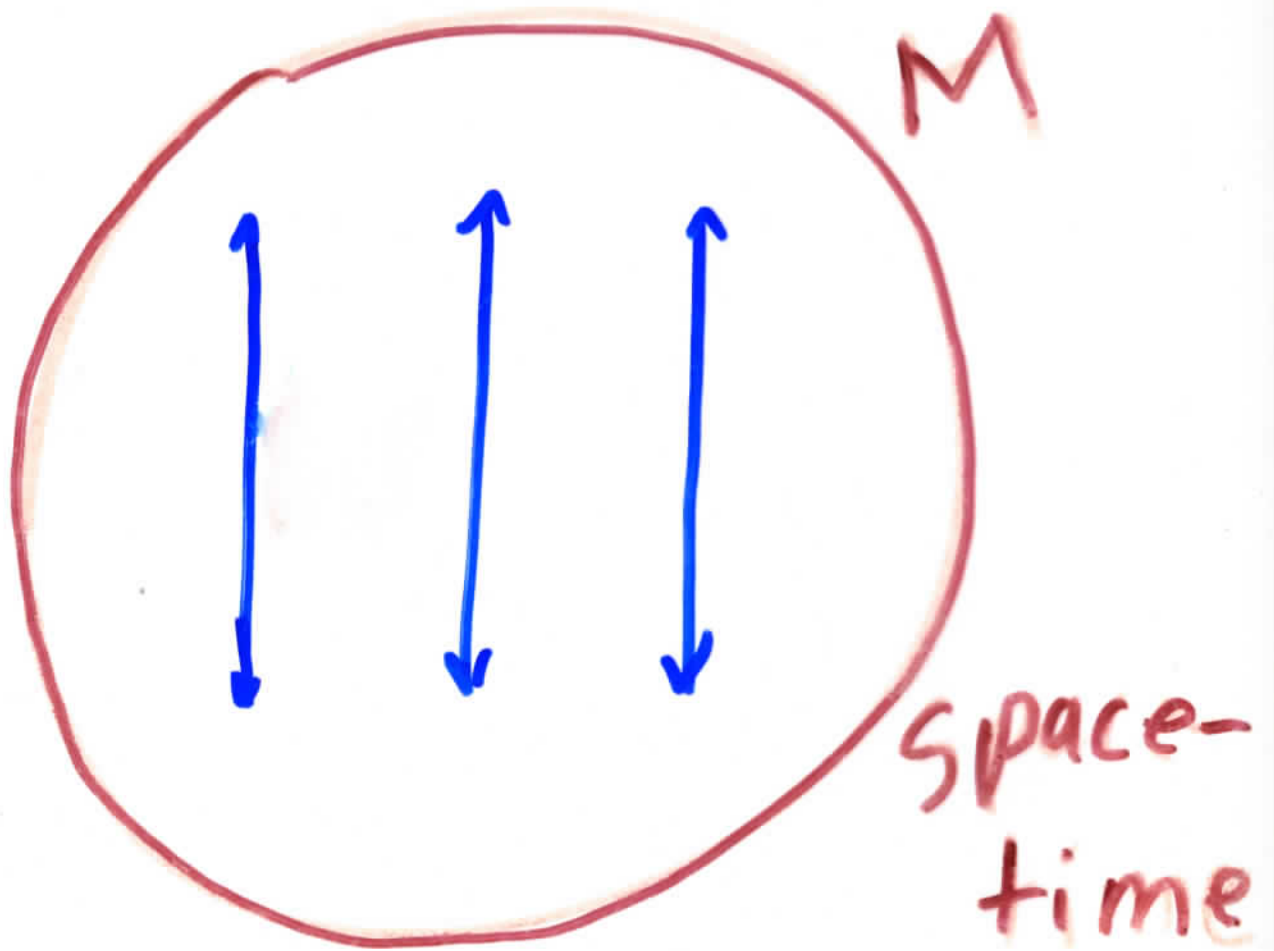
strongly causal and
compactness of
causal curves



physical
predictability!

stationary:
timelike Killing field K

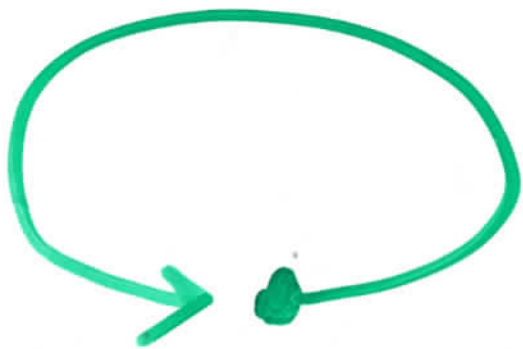




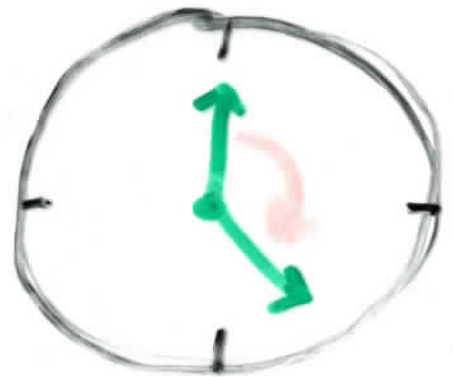
space of stationary orbits

Causal curiosity:

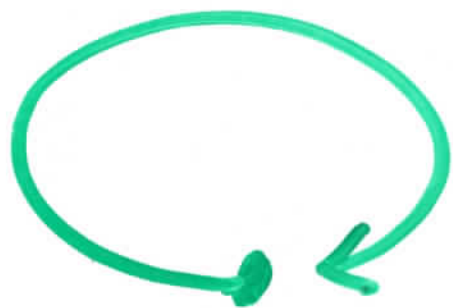
photon track:



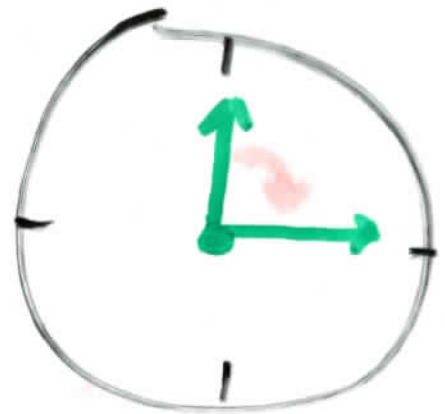
track-length
= L



T_1

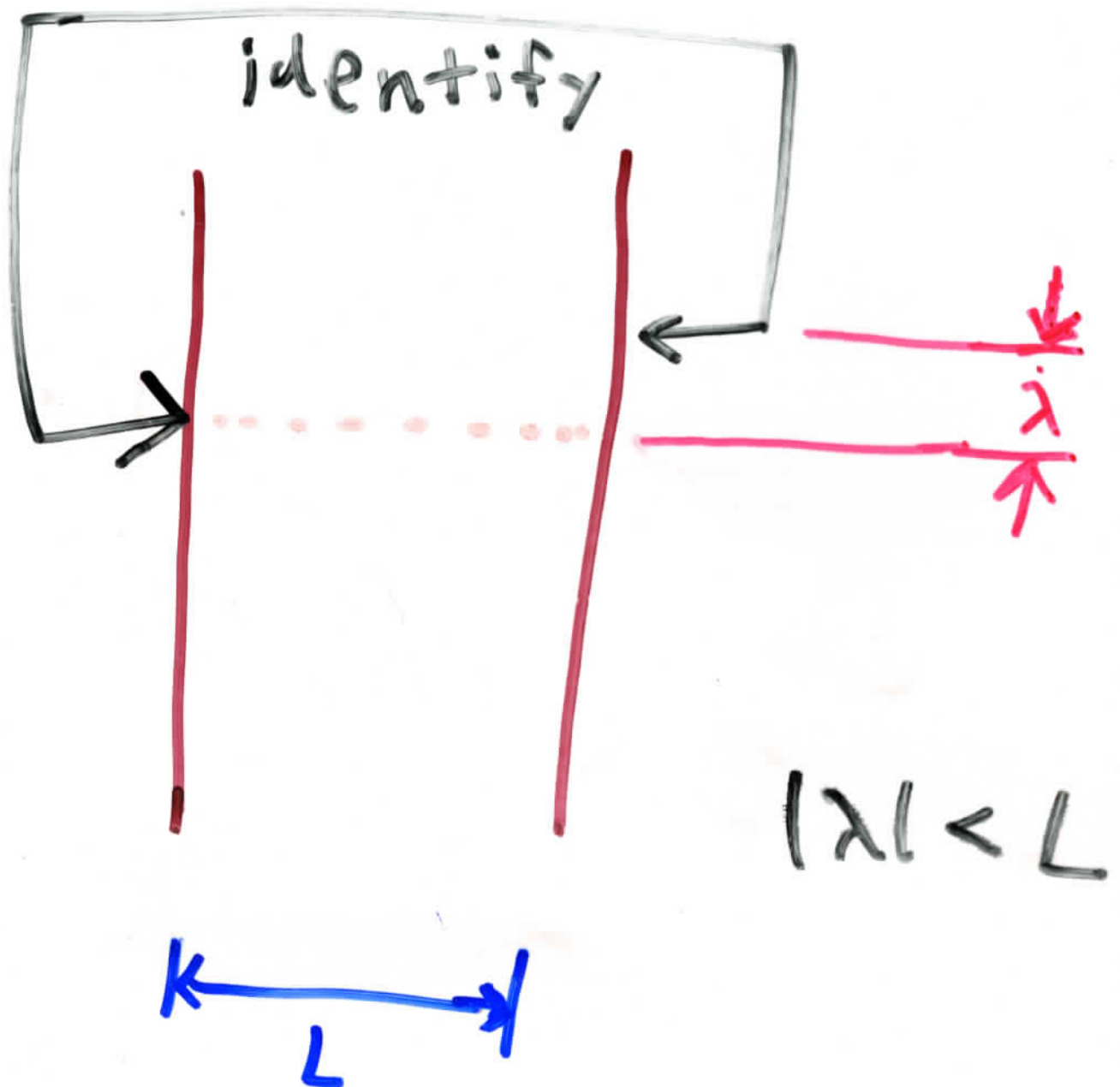


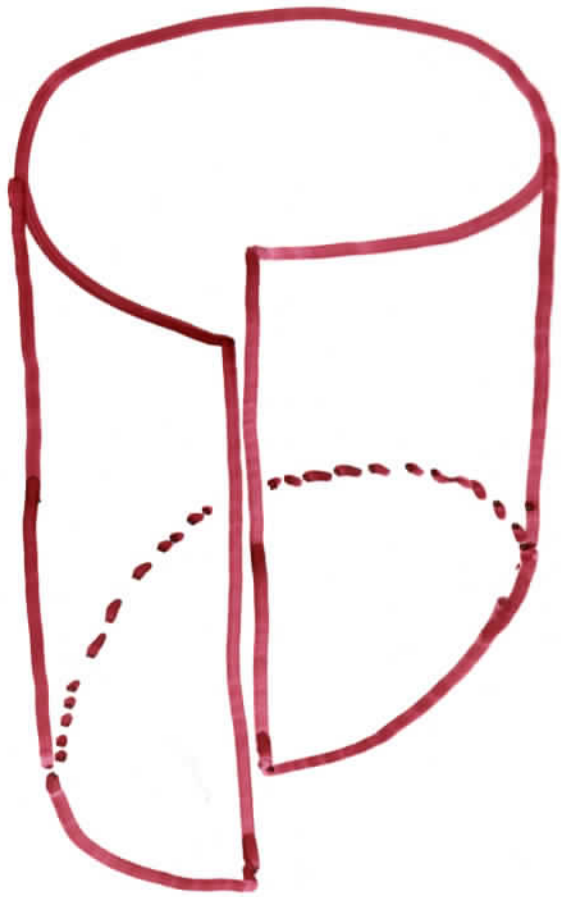
$T_1 \neq T_2$



T_2

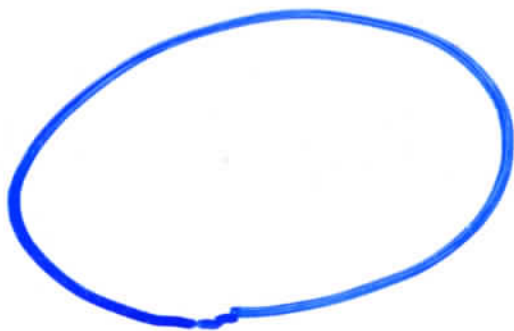
1+1 Minkowski cylinder





M

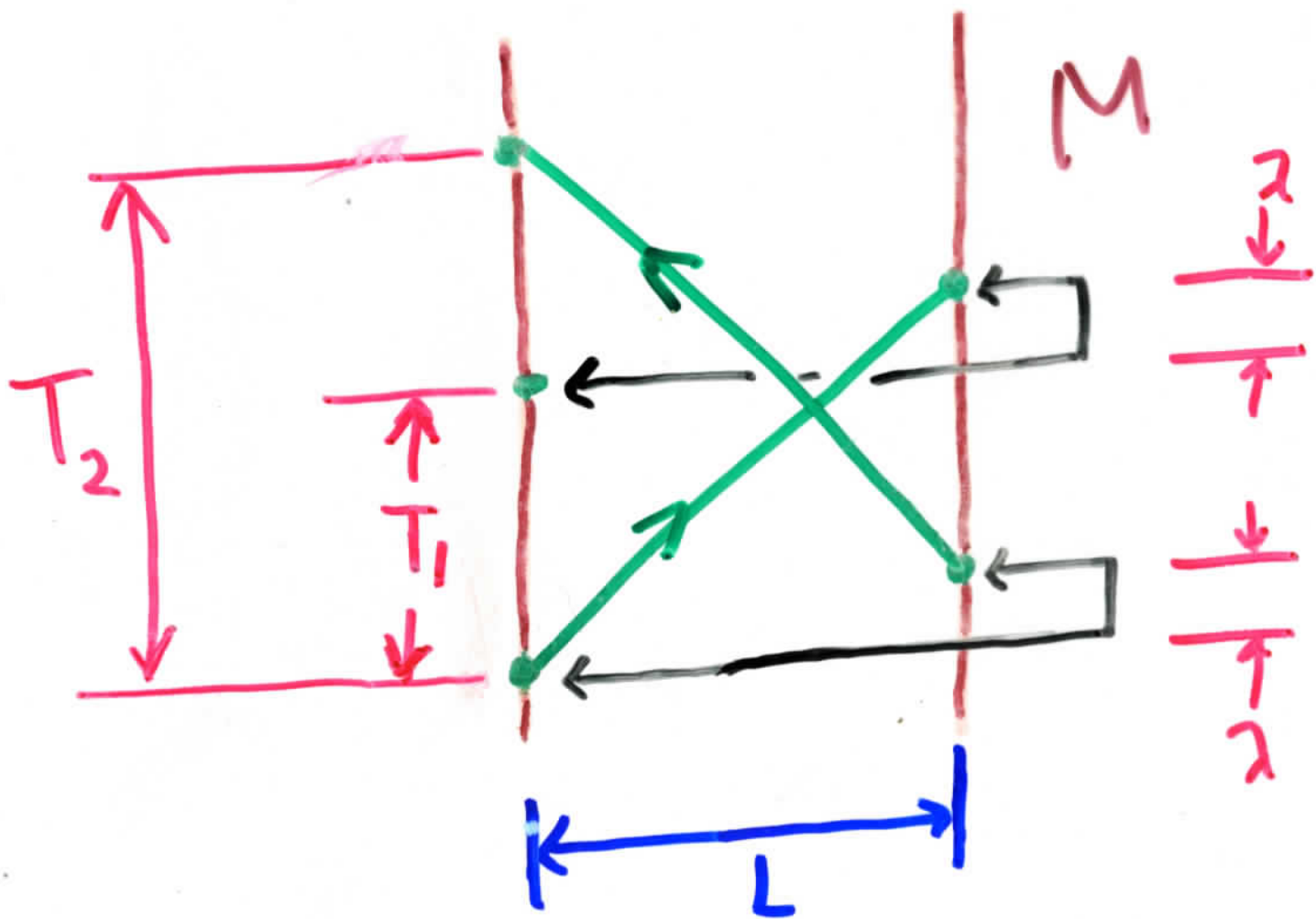
$\downarrow \pi$



Q

$$T_1 = L - \lambda$$

$$T_2 = L + \lambda$$



all measured by one observer

COSMIC STRING

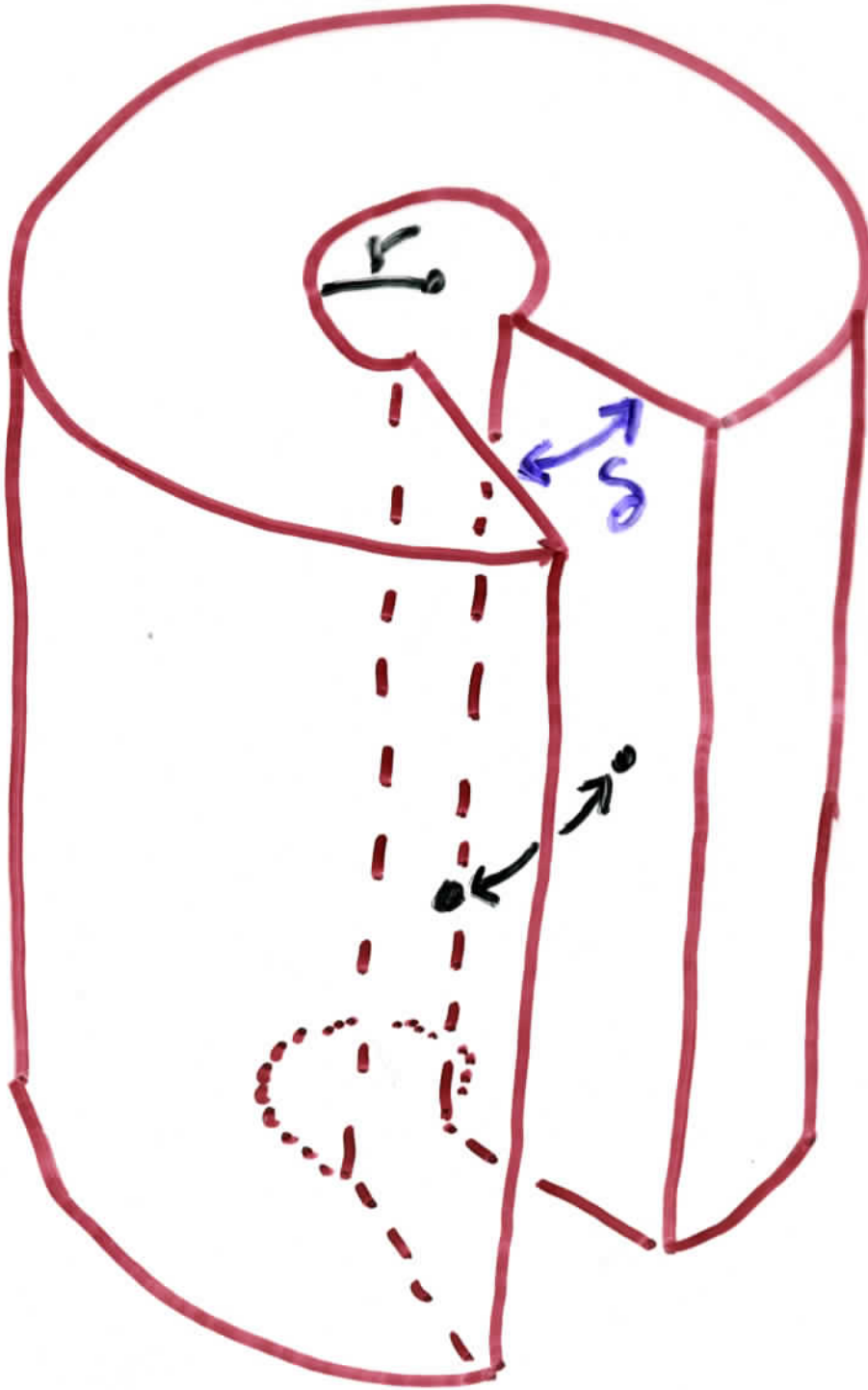
$3+0$



2+1 slice

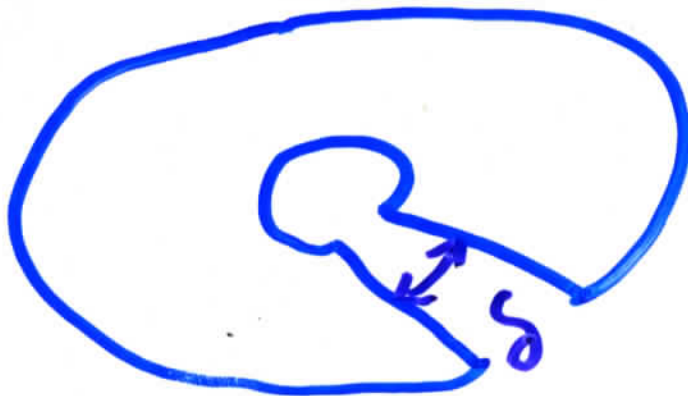


M



Ordinary
cosmic
string

Q

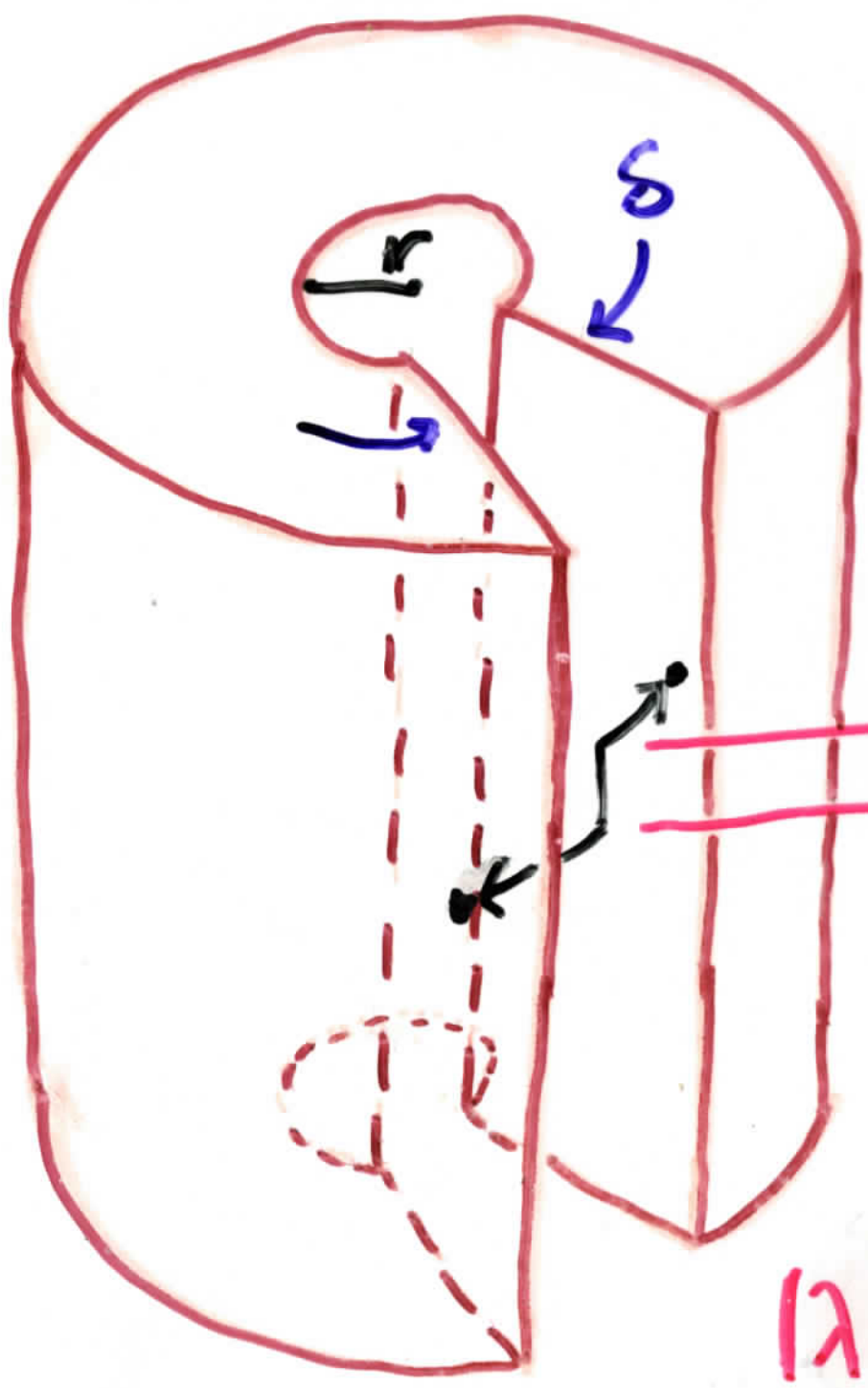


2+1 slice



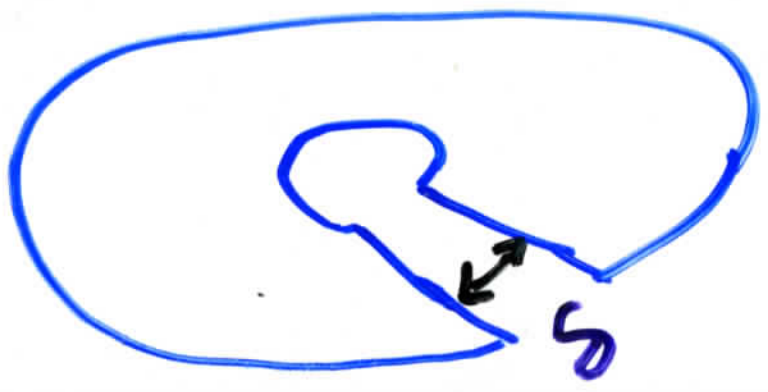
time-shifted
cosmic
string

M

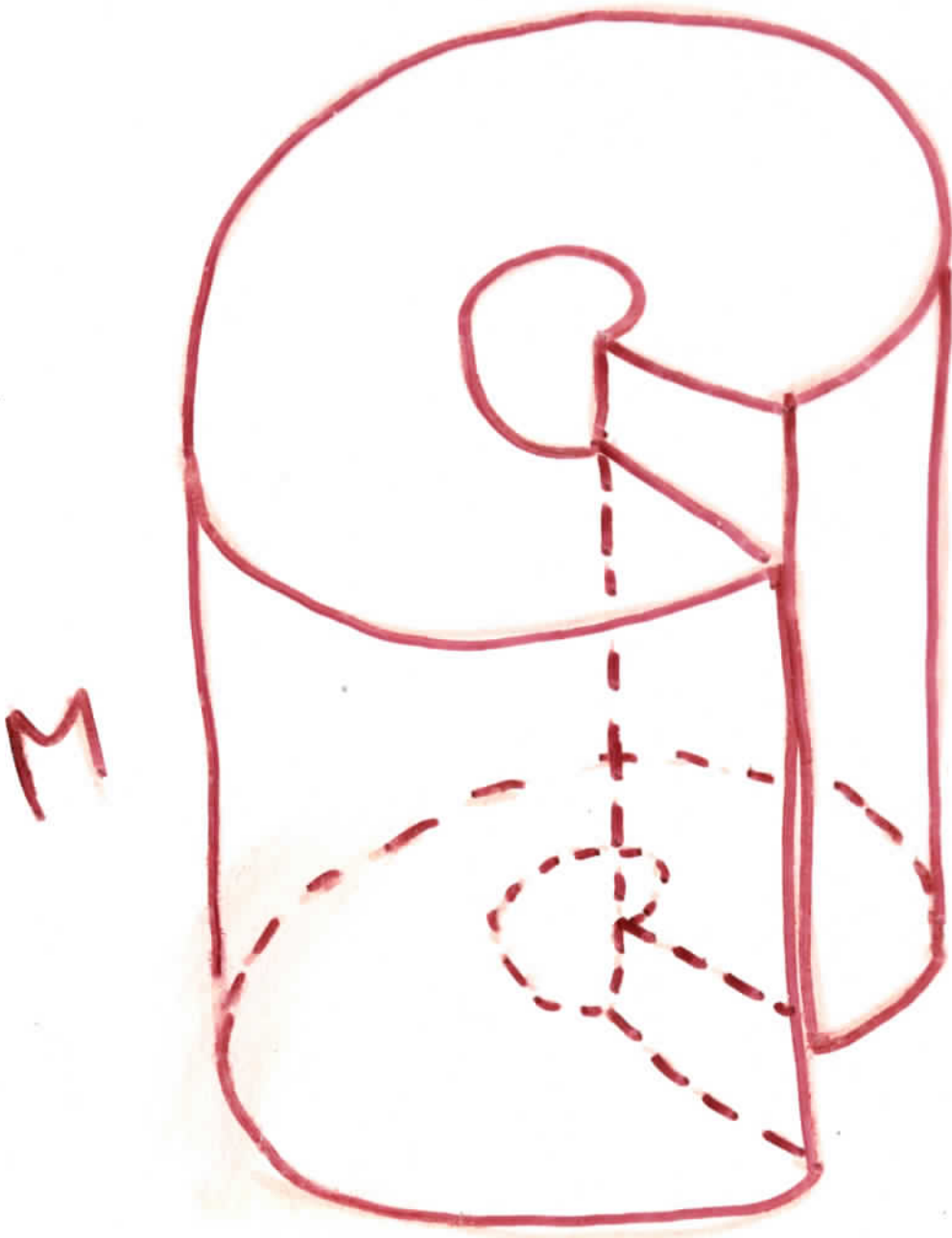


$|\lambda| < r(2\pi - \delta)$

Q

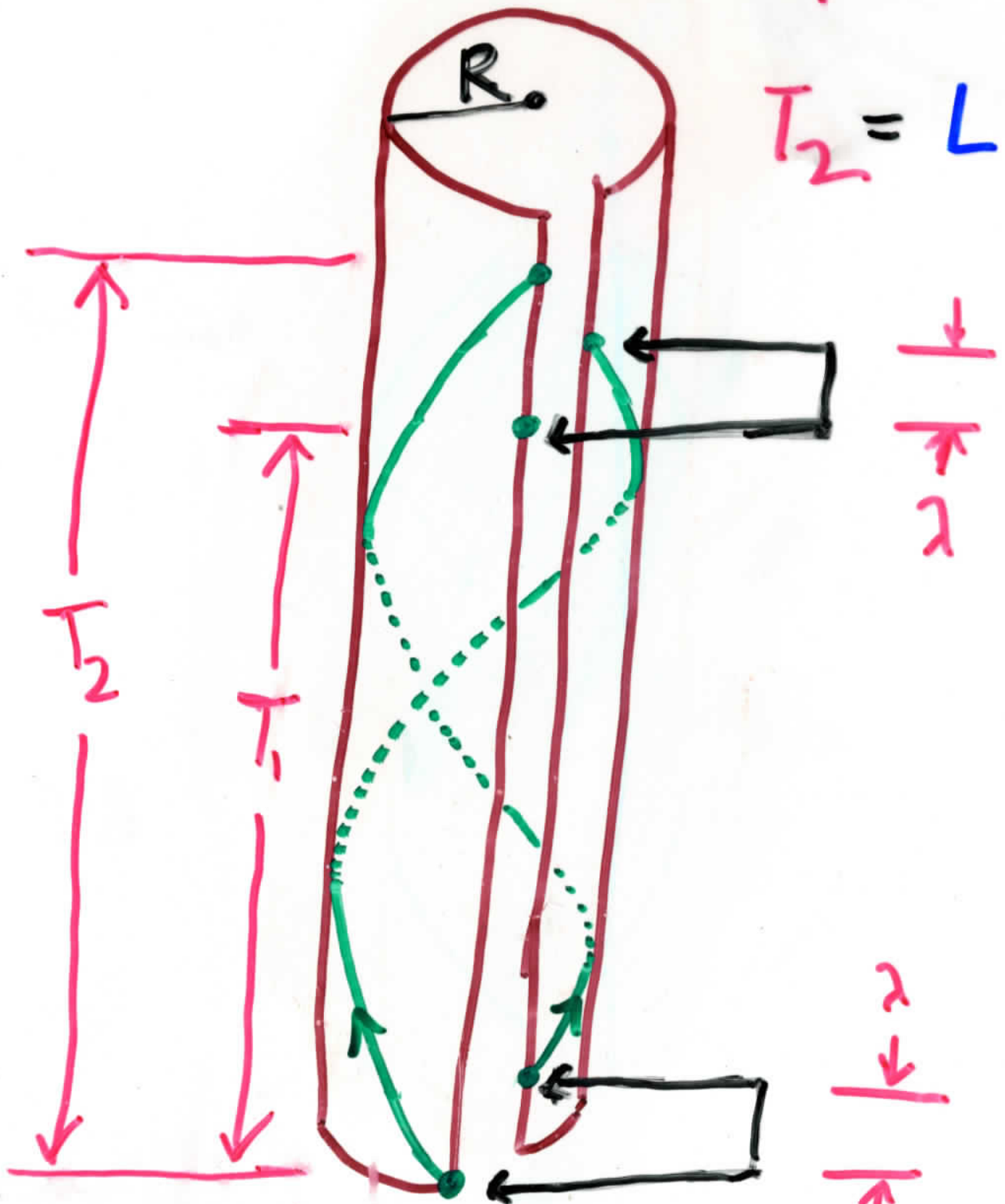


2+1



$$T_1 = L - \lambda$$

$$T_2 = L + \lambda$$



measured
by one
observer

$$L = R(2\pi - \delta)$$

causal

curiosity:

M ← GLOBALLY HYPERBOLIC

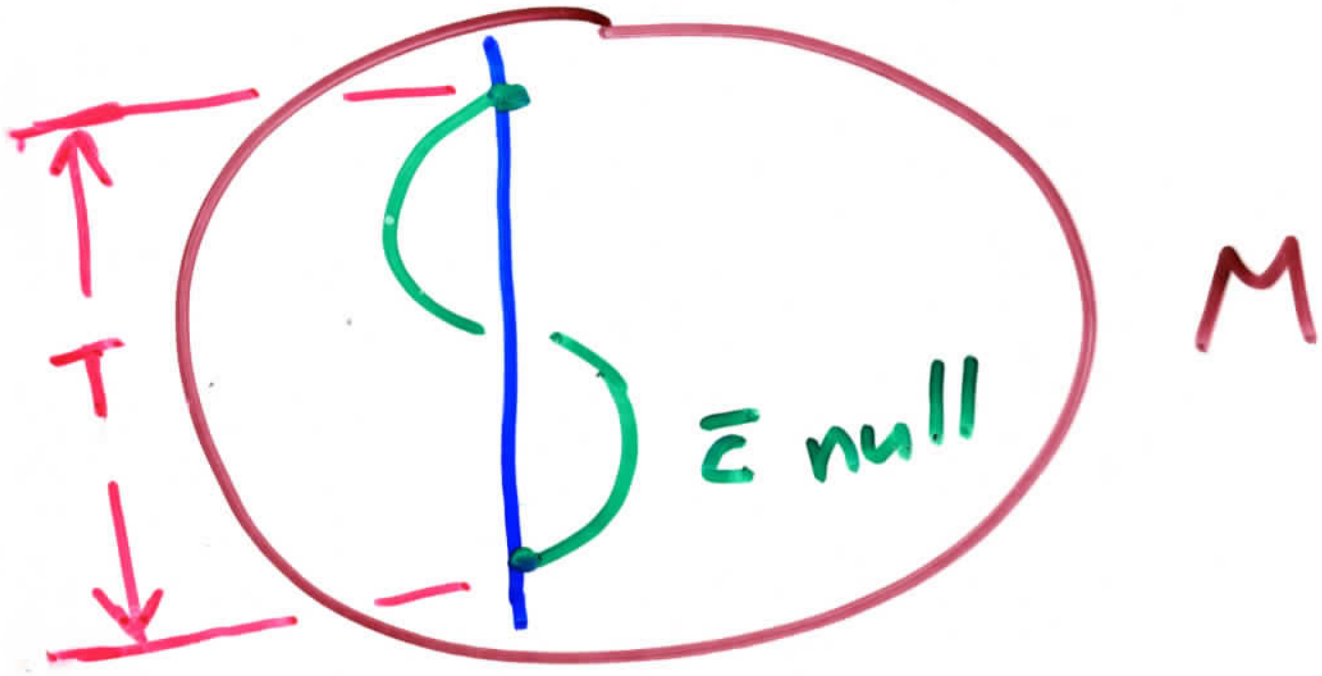
↓ quotient

M' assume:
strongly causal

↑ GLOBALLY HYPERBOLIC?

FUNDAMENTAL

COCYCLE β

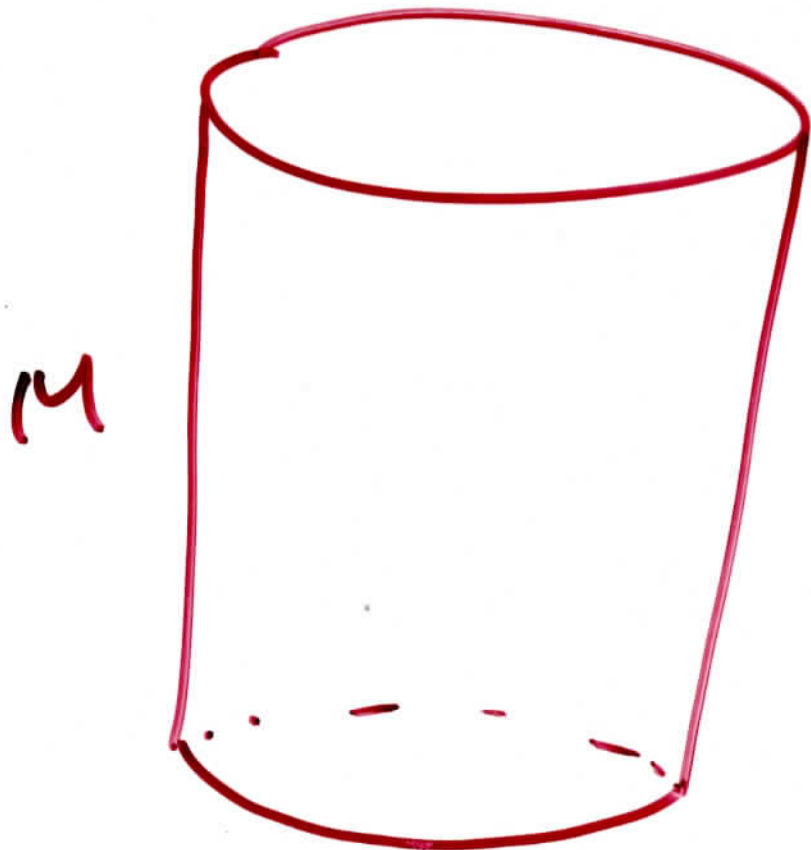


$$\beta(c) = T - L$$

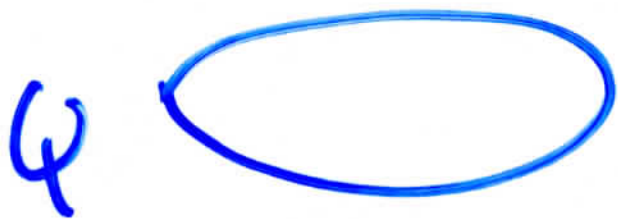


$$L = L(c)$$

metric g on M



$\downarrow \pi$



Riem. metric h_Q

metric g on M

$$(d\tau)K = 1$$

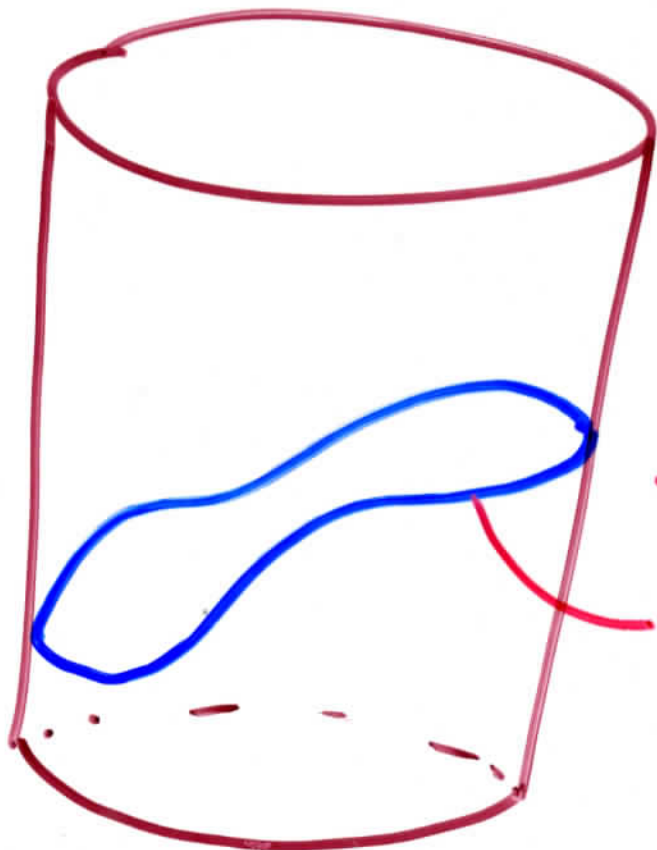
lots of freedom

γ

$\gamma=0$

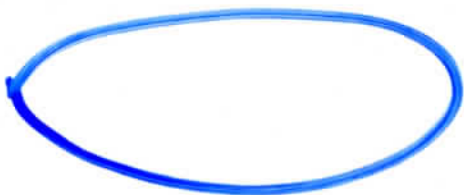
\mathbb{R}

M



π

Q



Riem. metric h_Q

metric g on M

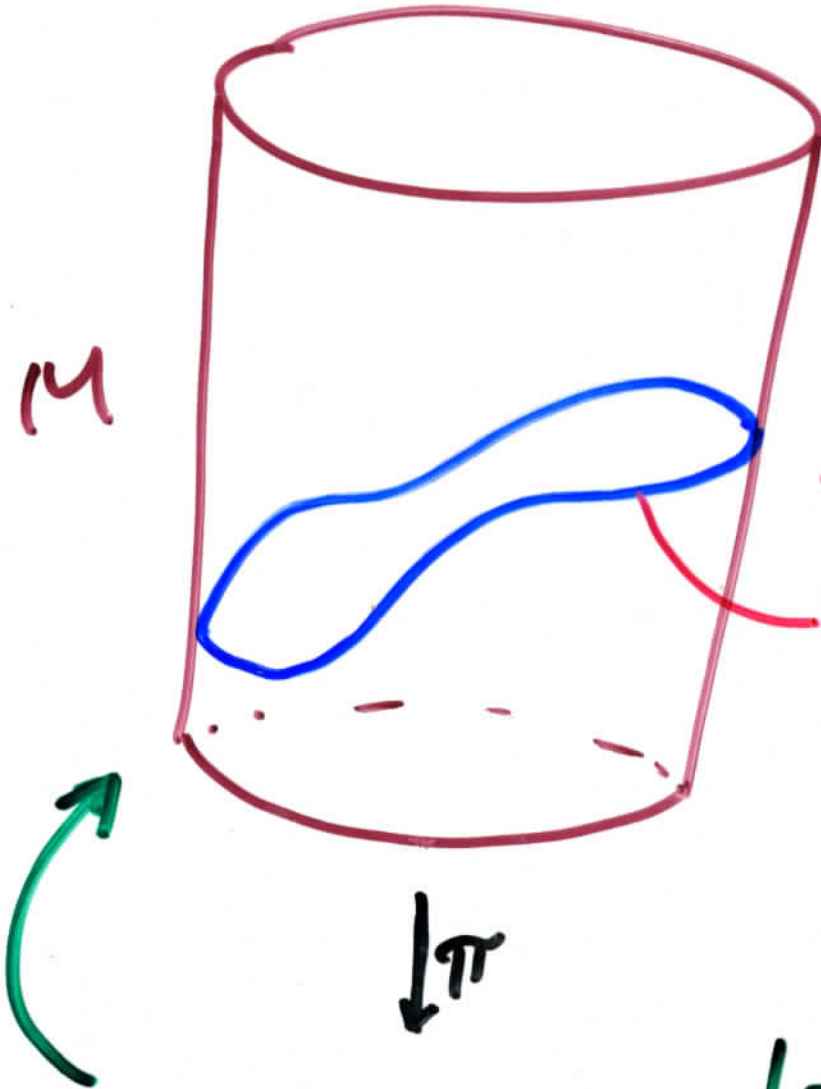
$$(d\tau)K = 1$$

lots of freedom

γ

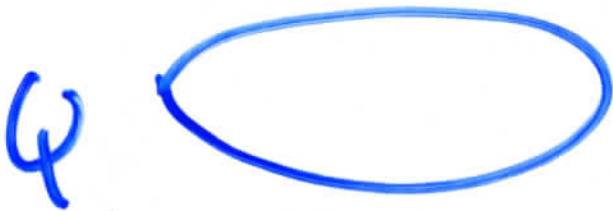
$\gamma=0$

\mathbb{R}



1-form ω^γ

on Q

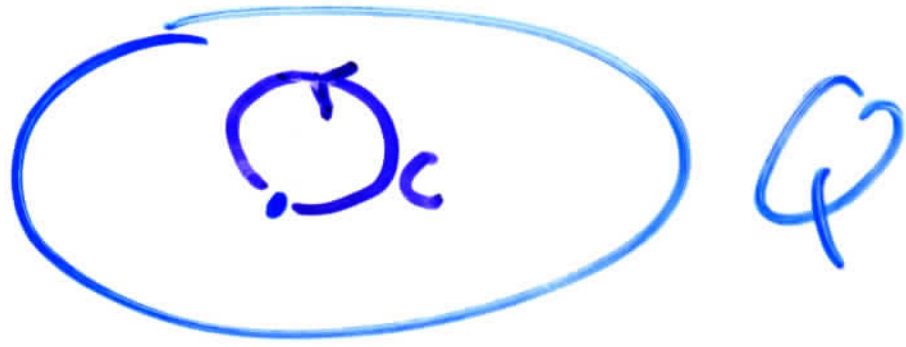


Riem. metric h_Q

(drift form)

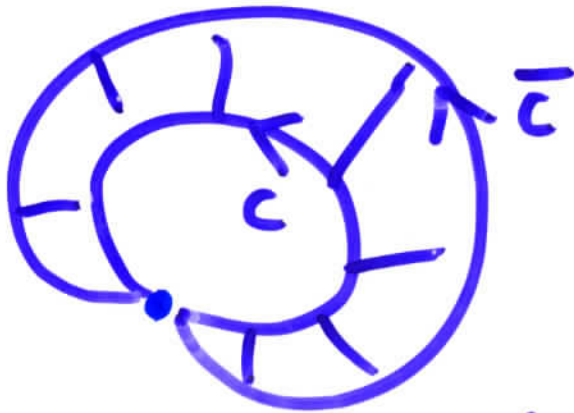
$$g = -|K|^2 (d\tau + \omega)^\gamma + h_Q$$

static : $d\omega = 0$



$$B(c) = \int_c \omega$$

static: $(d\omega = 0)$



c, \bar{c}

homotopic \Rightarrow

$$B(c) = B(\bar{c})$$

$$B = [\omega] \in H_{dR}^1(Q)$$



$$Q + B + |K|^2$$

Uniquely
determine



M

B, \bar{B} agree on

null-homotopic c



M, \bar{M} share cover



Weight of β

$$wt(\beta) = \sup_c \frac{\beta(c)}{L(c)}$$

RESULTS

$wt(\beta) > 1$: chronologically vicious

$wt(\beta) < 1$: strongly causal

M globally hyperbolic \iff
 \mathcal{G} complete

$$Wt(\beta) = 1:$$

RESULTS



$\beta(c) = L(c)$: chron.,
not causal

$$\frac{\beta(c_n)}{L(c_n)} \rightarrow 1$$

for some
 $\{c_n\}$

$$L(c_n) - \beta(c_n) \rightarrow 0:$$



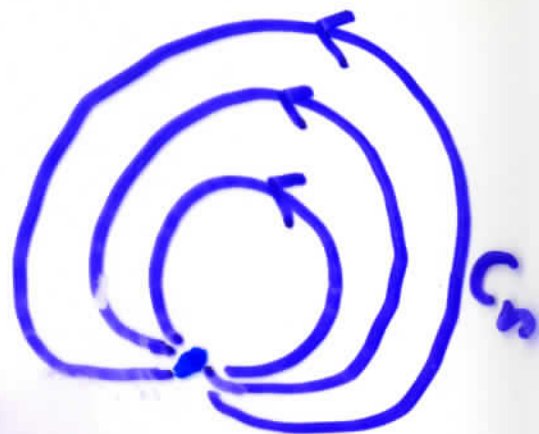
causal,
not strongly
causal

RESULTS

$$\text{wt}(\beta) = 1:$$

$$\frac{\beta(c_n)}{L(c_n)} \longrightarrow 1$$

for all $\{c_n\}$



$$L(c_n) - \beta(c_n) > \delta$$

for some $\{c_n\}$

$$L(c_n) - \beta(c_n) < A:$$

Strongly causal,
not glob. hyper.

$$L(c_n) - \beta(c_n) \longrightarrow \infty:$$

for all $\{c_n\}$

glob. hyper. \iff
 \mathcal{C} complete

RESULTS

M glob. hyper.

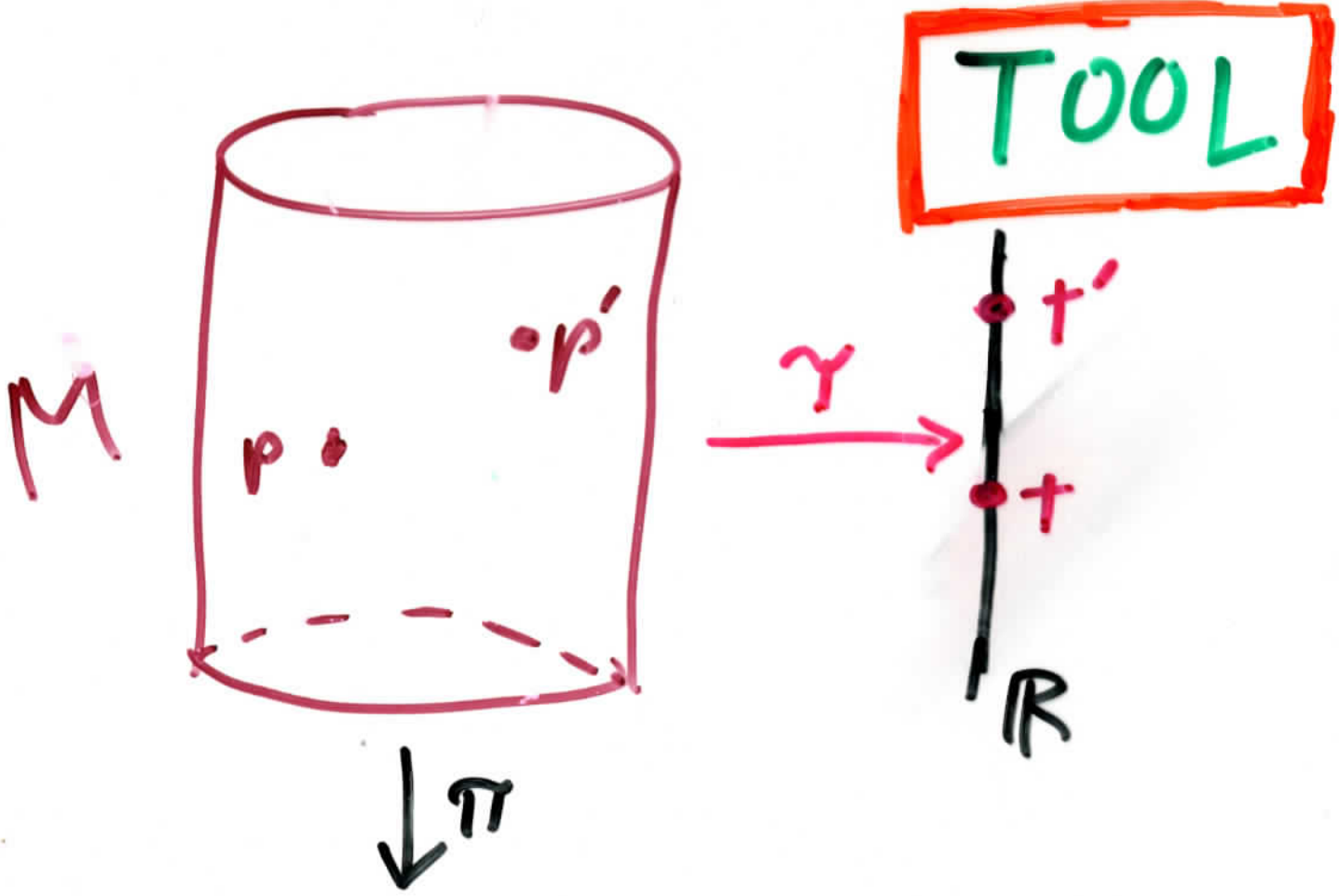


M'

M, M' stationary

$wt(\beta_{M'}) < 1 \Rightarrow$

M' glob. hyper.

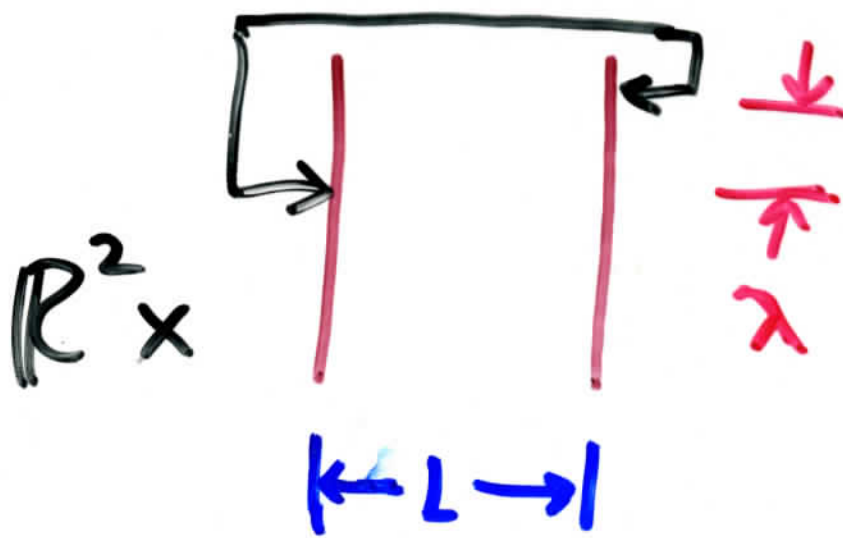


$$p \ll p' \iff$$

$$t' - t > \inf_c (L(c) - \int_c \omega)$$

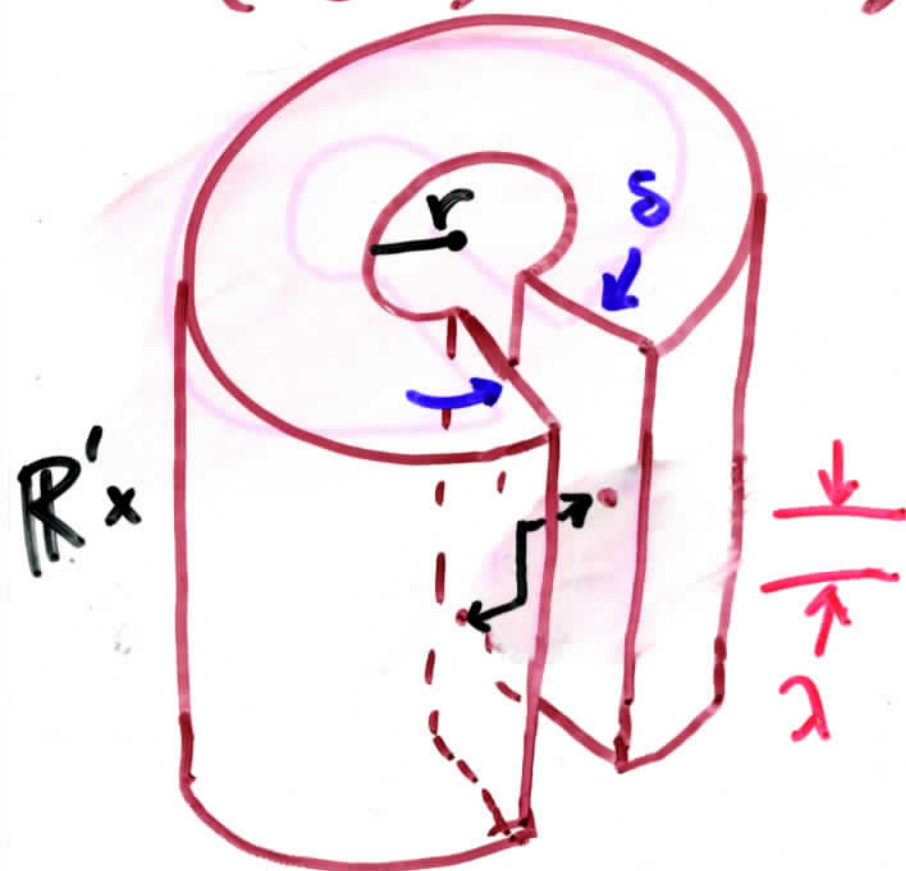
quasi-distance

3 + 1 Minkow. cylinder:



$$\text{wt}(\beta) = \frac{|\lambda|}{L}$$

cosmic string:



$$\text{wt}(\beta) = \frac{|\lambda|}{r(2\pi - \delta)}$$

Kerr

$$r > 2m$$

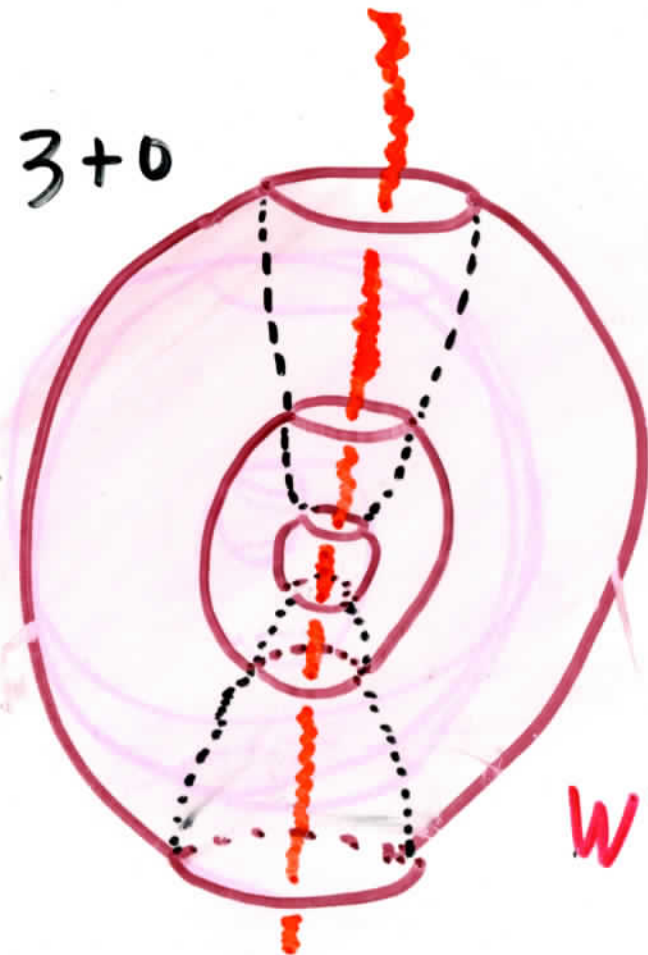
$$a < m$$

Spinning
black
hole

$$wt(\mathcal{B}) = 1$$

Kerr cosmic string

3+0



excise
polar caps

angle $\rightarrow 0$
for $r \rightarrow \infty$

$$wt(\mathcal{B}) = \text{ad lib.}$$

Kerr

$$\omega = \frac{2 m r a \sin^2 \theta}{r^2 - 2 m r + a^2 \cos^2 \theta} d\phi$$

ϕ = longitude

θ = co-latitude

arXiv: 1412.7742